#### 6 Finite Element Computation Scheme of Elasticity Problems

- 6.1 Weak form for general elasticity problems
- 6.2 Finite element method for solving elasticity problems

- 6.3 Global assembly from high dimensional elements
- 6.4 Treatments on boundary conditions
- **6.5 Exercises**

#### 6.3 Global assembly from high dimensional elements

- The global assembly of two-dimensional or three-dimensional elements is exactly the same as that of the previous one-dimensional problem according to the element location vector.
- For the convenience of introduction, a typical two-dimensional example is given to explain the global assembly process and computation procedure.









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6.3 Global	l assembly fron	ı high dimensiona	l elements	
$\mathbf{K}^{e} = \begin{bmatrix} \mathbf{K}_{11}^{e} & \mathbf{K}_{12}^{e} \\ \mathbf{K}_{21}^{e} & \mathbf{K}_{22}^{e} \\ \mathbf{K}_{31}^{e} & \mathbf{K}_{32}^{e} \\ \mathbf{K}_{41}^{e} & \mathbf{K}_{42}^{e} \end{bmatrix}$	$ \begin{array}{c} \mathbf{K}_{13}^{e} & \mathbf{K}_{14}^{e} \\ \mathbf{K}_{23}^{e} & \mathbf{K}_{24}^{e} \\ \mathbf{K}_{33}^{e} & \mathbf{K}_{34}^{e} \\ \mathbf{K}_{43}^{e} & \mathbf{K}_{44}^{e} \end{array} $			
=	240384.60 - 42735.03 758546.95 48076.92 491480.77 Symmetric	- 48076.92 - 245726.48 - 630341.82 - 240384.60 - 240384.60 202991.45 758546.95	-240384.60 -202991.45 -379273.48 -48076.92 -48076.92 -245726.48 -251 <u>068.35</u> 240384.60 -42735.03 758546.95 48076.92 491480.77	48076.92 251068.35 240384.60 - 379273.48 - 48076.92 - 630341.82 - 240384.60 758546.95
• In this numer the sam values	way, the stiffno ical integration ne computatior of stiffness ma	ess coefficient can further, other stif process. Through trices of each elen	be obtained by fness coefficients the above steps, nent can be obtain	have the 1919



			in mgn unn				
$\lambda^{1}$	= {2	1 3	4} <sup><i>T</i></sup>	$\lambda^2 = \{4$	3 5	6} <sup>1</sup>	τ
Node	1	2	3	4	5	6	
number	$1 \left[ \mathbf{K} \right]$	$K_{21}^{1}$	$K_{23}^{1}$	$\mathbf{K}_{24}^1$		٦٢	$\hat{\mathbf{u}}_1$ $[\mathbf{f}_1]$
	2 K	$\mathbf{K}_{12}^{1}$	$\mathbf{K}_{13}^{1}$	$\mathbf{K}_{14}^{1}$			$\hat{\mathbf{u}}_2$ $\mathbf{f}_2$
	3 K	${\bf K}_{31}^1$	$(\mathbf{K}_{33}^1 + \mathbf{K}_{22}^2)$	$(\mathbf{K}_{34}^1 + \mathbf{K}_{21}^2)$	$K_{23}^{2}$	$\mathbf{K}_{24}^{2}$	$\hat{\mathbf{u}}_3$ $\mathbf{f}_3$
	4 <b>K</b>	$K_{41}^1$	$\left(\overline{\mathbf{K}_{43}^{1}}+\mathbf{K}_{12}^{2}\right)$	$(\mathbf{K}_{44}^{1} + \mathbf{K}_{11}^{2})$	$K_{13}^{2}$	$\mathbf{K}_{14}^2$	$\hat{\mathbf{u}}_4 = \mathbf{f}_4$
	5		$K_{32}^{2}$	$K_{31}^{2}$	$K_{33}^{2}$	$\mathbf{K}_{34}^{2}$	$\hat{\mathbf{u}}_{5}$ $\mathbf{f}_{5}$
	6		$K_{42}^{2}$	$K_{41}^{2}$	$K_{43}^{2}$	$\mathbf{K}_{44}^2$	$\hat{\mathbf{u}}_{6}$ $[\mathbf{f}_{6}]$
	-					_	





### 6.4 Treatments on boundary conditions

#### (1) Displacement boundary conditions

• With the above finite element form, it is very simple to impose the displacement boundary conditions since the parameters are now all physical values. That is they obey the property







<b>5.3 Globa</b> $\mathbf{F} = \begin{bmatrix} \mathbf{K}_{11}^{e} & \mathbf{K}_{12}^{e} \\ \mathbf{K}_{21}^{e} & \mathbf{K}_{22}^{e} \\ \mathbf{K}_{31}^{e} & \mathbf{K}_{32}^{e} \\ \mathbf{K}_{41}^{e} & \mathbf{K}_{42}^{e} \end{bmatrix}$		bly fron	n high dii	mensiona	l elemen	ts	
=	240384.60 758546.95 Symmetric	- 42735.03 48076.92 491480.77	- 48076.92 - 630341.82 - 240384.60 758546.95	- 245726.48 - 240384.60 202991.45 - <u>48026.92</u> 491480.77	-240384.60 -379273.48 -48076.92 251068.35 240384.60 758546.95	- 202991.45 - 48076.92 - 245726.48 240384.60 - 42735.03 48076.92 491480.77	48076.92 251068.35 240384.60 - 379273.48 - 48076.92 - 630341.82 - 240384.60 758546.95
<ul> <li>In this numer the same values</li> </ul>	way, the rical interior of stiff	e stiffne egration putation	ess coeffi , further, 1 process. trices of	cient can other stif Through each elen	be obtain fness coe the above nent can	ned by fficients ve steps, t be obtain	have the







#### 6.4 Treatments on boundary conditions

· The displacement functions of each element can be obtained by the combination of shape function and node displacements

Element 1  

$$\hat{u}^{1}(\xi,\eta) = \sum_{\sigma=1}^{4} N_{\sigma}(\xi,\eta) \hat{u}_{\sigma}^{1} = N_{1} \cdot 0 + N_{2} \cdot 0 + N_{3} \cdot 0 + N_{3}$$

$$\hat{v}^{1}(\xi,\eta) = \sum_{\sigma=1}^{4} N_{\sigma}(\xi,\eta) \hat{v}_{\sigma}^{1} = N_{1} \cdot 0 + N_{2} \cdot 0 + N_{3} \cdot (-0)$$

$$= -0.1648 (1+\xi)(1+\eta)$$



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# $(x_1, y_1) = (0, -10), (x_2, y_2) = (0, 0)$ $\hat{u}^{1}(\xi,\eta) = \sum_{a=1}^{4} N_{a}(\xi,\eta) \hat{u}_{a}^{1} = N_{1} \cdot 0 + N_{2} \cdot 0 + N_{3} \cdot 0 + N_{4} \cdot 0 = 0$ $\hat{v}^{1}\left(\xi,\eta\right) = \sum_{i=1}^{4} N_{\sigma}\left(\xi,\eta\right) \hat{v}_{\sigma}^{1} = N_{1} \cdot 0 + N_{2} \cdot 0 + N_{3} \cdot \left(-0.6592\right) + N_{4} \cdot 0$ $= -0.1648(1+\xi)(1+\eta) = -0.004395 x(y+10)$

#### 6.4 Treatments on boundary conditions

• The displacement functions of each element can be obtained by the combination of shape function and node displacements

#### Element 2

$$\hat{u}^{2}(\xi,\eta) = \sum_{a=1}^{4} N_{a}(\xi,\eta) \hat{u}_{a}^{2} = N_{1} \cdot 0 + N_{2} \cdot 0 + N_{3} \cdot 0 + N_{4} \cdot 0 = 0$$
$$\hat{v}^{2}(\xi,\eta) = \sum_{a=1}^{4} N_{a}(\xi,\eta) \hat{v}_{a}^{2} = N_{1} \cdot 0 + N_{2} \cdot (-0.6592) + N_{3} \cdot 0 + N_{4} \cdot 0$$
$$= -0.1648 (1 - \xi)(1 + \eta)$$

#### 6.4 Treatments on boundary conditions

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· Once the displacement are known, the approximation for stresses on element may be computed by

$$\hat{\boldsymbol{\sigma}} = \sum_{b=1}^{4} \mathbf{DS}(N_b) \hat{\mathbf{u}}_b^e = \sum_{b=1}^{4} \mathbf{DB}_b \hat{\mathbf{u}}_b^e$$

• In the computation process, the derivatives of each shape function is needed, which is the same as the way for  $N_3$  above. The final derivative results of other shape functions are given below





## 6.4 Treatments on boundary conditions 6.4 Treatments on boundary conditions The stress functions of each element can be obtained by the combination of derivatives of shape functions and node displacements displacements **Element 2** $\hat{\sigma}^2 = \{\sigma_x^2 \quad \sigma_y^2 \quad \sigma_z^2 \quad \tau_{xy}^2\}^{T} = \sum_{b=1}^4 \mathbf{DS}(N_b) \hat{\mathbf{u}}_b^2 = \sum_{b=1}^4 \mathbf{DB}_b \hat{\mathbf{u}}_b^2$ $= \mathbf{D} \Big( \mathbf{B}_1 \hat{\mathbf{u}}_1^2 + \mathbf{B}_2 \hat{\mathbf{u}}_2^2 + \mathbf{B}_3 \hat{\mathbf{u}}_3^2 + \mathbf{B}_4 \hat{\mathbf{u}}_4^2 \Big) = \mathbf{D} \Big( \mathbf{B}_1 \mathbf{0} + \mathbf{B}_2 \hat{\mathbf{u}}_2^2 + \mathbf{B}_3 \mathbf{0} + \mathbf{B}_4 \mathbf{0} \Big) = \mathbf{D} \mathbf{B}_2 \hat{\mathbf{u}}_2^2$ $$\begin{split} &= \mathbf{D}(\mathbf{B}, \vec{\mathbf{u}}_1^+ + \mathbf{B}_2 \vec{\mathbf{u}}_2^+ + \mathbf{B}_1 \vec{\mathbf{u}}_1^+) = \mathbf{D}(\mathbf{B}, \mathbf{0} + \mathbf{B}_1 \vec{\mathbf{u}}_2^+ + \mathbf{B}_1 \mathbf{0} + \mathbf{B}_1 \mathbf{0}) = \mathbf{D}(\mathbf{B}) \\ &= \frac{E}{d} \begin{bmatrix} (1-\nu) & \nu & \nu & 0 \\ \nu & (1-\nu) & \nu & 0 \\ 0 & 0 & 0 & (1-2\nu)/2 \end{bmatrix} \begin{bmatrix} \frac{\partial N_z}{\partial \mathbf{v}} & 0 \\ 0 & \frac{\partial N_z}{\partial \mathbf{v}} \end{bmatrix} \begin{bmatrix} \hat{\mu}_z^2 \\ \hat{\mu}_z^2 \\ \hat{\nu}_z^2 \end{bmatrix} \\ &= \frac{E}{d} \begin{bmatrix} (1-\nu) & \nu & \nu & 0 \\ \nu & (1-\nu) & \nu & 0 \\ \nu & \nu & (1-\nu) & \nu & 0 \\ 0 & 0 & 0 & (1-2\nu)/2 \end{bmatrix} \begin{bmatrix} -\frac{1}{30}(1+\eta) & 0 \\ 0 & \frac{1}{20}(1-\xi) \\ 0 & \frac{1}{20}(1-\xi) \end{bmatrix} \\ &= \frac{1}{20} \begin{bmatrix} 1-\frac{1}{20} \\ 1-\frac$$ - 0.6592 37 37 37



#### 6.4 Treatments on boundary conditions



· Through the solutions of the above stress functions, the stress values at the integration points can be obtained, by substituting the coordinates of the integration points of the specified element. For example, the stress at integral point  $(\xi, \eta) = \left(-\frac{1}{\sqrt{3}}, \frac{1}{\sqrt{3}}\right)$ of element 1 is  $\begin{bmatrix} 1 \\ 1 \\ 1 \end{bmatrix}$ 

$$\hat{\boldsymbol{\sigma}}^{1} = \left\{ \sigma_{z}^{1} \quad \sigma_{y}^{1} \quad \sigma_{z}^{1} \quad r_{zy}^{1} \right\}^{T} = \left\{ \begin{matrix} -19015.38(1+\xi) \\ -44369.23(1+\xi) \\ -19015.38(1+\xi) \\ -19015.38(1+\xi) \\ -8451.28(1+\eta) \end{matrix} \right\} = \left\{ \begin{matrix} -19015.38(1-\frac{1}{\sqrt{3}}) \\ -19015.38(1-\frac{1}{\sqrt{3}}) \\ -8451.28(1+\frac{1}{\sqrt{3}}) \end{matrix} \right\} = \left\{ \begin{matrix} -8036.84 \\ -18752.63 \\ -8036.84 \\ -13330.63 \end{matrix} \right\}$$

## 6.4 Treatments on boundary conditions

· Furthermore, by substituting the node coordinates of the specified element into the solution of the above stress function, the stress values at the nodes can be obtained. For example, the stresses at node 2 of  $(\xi, \eta) = (-1, 1)$  of element 1 are as follows  $\hat{\boldsymbol{\sigma}}^{1} = \left\{ \sigma_{z}^{1} \quad \sigma_{y}^{1} \quad \sigma_{z}^{1} \quad \tau_{zy}^{1} \right\}^{T} = \left\{ \begin{matrix} -19015.38(1+\xi) \\ -44369.23(1+\xi) \\ -19015.38(1+\xi) \\ -8451.28(1+\eta) \end{matrix} \right\} = \left\{ \begin{matrix} -19015.38(1-1) \\ -44369.23(1-1) \\ -19015.38(1-1) \\ -8451.28(1+1) \end{matrix} \right\} = \left\{ \begin{matrix} 0 \\ 0 \\ -16902.56 \end{matrix} \right\}$ 40 40

