

### 3 Weak Form of Equivalent Integration



#### 3.1 Weak form of equivalent integration for differential equations

#### 3.2 Weak form of one-dimensional elasticity problems

#### 3.3 Finite element computation based on weak form

#### 3.4 Global assembly from one-dimensional elements

#### 3.5 Treatments on boundary conditions

#### 3.6 Exercises

### 3.3 Finite element computation based on weak form



#### ■ Finite element computation

- A more convenient method to construct the approximating functions  $\phi_n$  and  $\psi_m$  are obtained by dividing the domain to be analyzed into small regular shaped regions. For example, we can divide the one-dimensional region between  $a$  and  $b$  into a set of “ $M$ ” small finite segments by defining a set of  $N$  points  $x_i$  such that

$$x_1 = a, \quad x_i < x_{i+1} \quad \text{and} \quad x_N = b$$

For a one-dimensional problem we can let each increment define a *finite element* domain (or more simply, an *element*) and the set of points define the *nodes* (*finite element mesh* or *mesh*).

2  
2

### 3.3 Finite element computation based on weak form



- A simple set of continuous polynomial approximating functions

$$\phi_i = \begin{cases} 0, & x < x_{i-1} \\ \frac{x - x_{i-1}}{x_i - x_{i-1}}, & x_{i-1} \leq x \leq x_i \\ \frac{x_{i+1} - x}{x_{i+1} - x_i}, & x_i < x \leq x_{i+1} \\ 0, & x > x_{i+1} \end{cases}$$

$C_0$  function since only the function is continuous in  $x$ , whereas the first derivative is only piecewise continuous with the discontinuities located at the nodes.

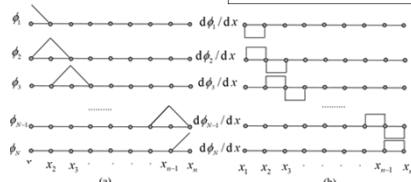


Figure 3.3 One-dimensional finite element approximation for  $\phi_i$ : (a) functions and (b) derivatives.

3  
3

### 3.3 Finite element computation based on weak form



- Integrals over each element in the weak form

$$\int_{\Omega} (\cdot) dx = \sum_{j=1}^M \int_{x_{j-1}}^{x_j} (\cdot) dx \equiv \sum_e \int_{\Omega_e} (\cdot) dx$$

Considering any interval  $[x_i, x_{i+1}]$ , we note that each interval is defined by the same two local functions  $N_1$  and  $N_2$ . We call these the *shape functions* (形函数) for the element.

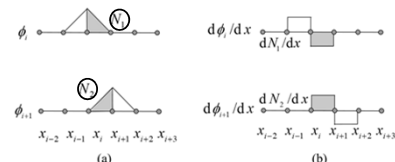


Figure 3.4 One-dimensional finite element shape functions: (a) functions and (b) derivatives.

4  
4

### 3.3 Finite element computation based on weak form



- Define local nodal coordinates  $x'$  on each element with  $x_1^e$  and  $x_2^e$ .

$$\begin{array}{c} e \\ \hline x_1^e \quad | \quad x' \quad | \quad x_2^e \\ \hline \end{array} \quad x' = x - x_1^e \quad h_e = x_2^e - x_1^e$$

- Local coordinate system

$$N_1(x') = 1 - \frac{x'}{h_e} \quad \text{and} \quad N_2(x') = \frac{x'}{h_e}$$

$$\frac{dN_1}{dx} = \frac{dN_1}{dx'} = -\frac{1}{h_e} \quad \text{and} \quad \frac{dN_2}{dx} = \frac{dN_2}{dx'} = \frac{1}{h_e}$$

- Displacements and arbitrary function on the element

$$\hat{u}^e = N_1(x') \hat{u}_1^e + N_2(x') \hat{u}_2^e$$

$$\hat{w}^e = N_1(x') \hat{w}_1^e + N_2(x') \hat{w}_2^e$$

5

### 3.3 Finite element computation based on weak form



- Weak form in global domain

$$G(w, u) = \int_{\Omega} w(x) b_x dx - \int_{\Omega} \frac{\partial w}{\partial x} E \frac{\partial u}{\partial x} dx + w \bar{f}_x \Big|_{\Gamma_f} = 0$$

- Weak form in each element

$$\hat{G}(\hat{w}, \hat{u}) = \hat{G}_e(\hat{w}, \hat{u}) - \hat{G}_f(\hat{w}, \hat{u}) - \hat{w}(x) \bar{f}_x \Big|_{\Gamma_f}$$

$$\hat{G}_e(\hat{w}, \hat{u}) = \sum_{e=1}^M \left[ \hat{w}_1^e \hat{w}_2^e \right]_0^{h_e} \left\{ \frac{dN_1}{dx'} \right\} E \left\{ \frac{dN_1}{dx'} \frac{dN_2}{dx'} \right\} dx' \left\{ \hat{u}_1^e \right\}$$

$$\hat{G}_f(\hat{w}, \hat{u}) = \sum_{e=1}^M \left[ \hat{w}_1^e \hat{w}_2^e \right]_0^{h_e} \left\{ \frac{N_1}{N_2} \right\} b_x dx'$$

6  
6

### 3.3 Finite element computation based on weak form

- Each element can be evaluated as element stiffness matrix and load matrix

$$\mathbf{K}^e = \int_0^{h_e} \begin{Bmatrix} \frac{dN_1}{dx'} \\ \frac{dN_2}{dx'} \end{Bmatrix} E \begin{bmatrix} \frac{dN_1}{dx'} & \frac{dN_2}{dx'} \end{bmatrix} dx' = \begin{bmatrix} K_{11}^e & K_{12}^e \\ K_{21}^e & K_{22}^e \end{bmatrix}$$

$$\mathbf{f}^e = \int_0^{h_e} \begin{Bmatrix} N_1 \\ N_2 \end{Bmatrix} b_s dx' = \begin{Bmatrix} f_1^e \\ f_2^e \end{Bmatrix}$$

- Element stiffness matrix and load matrix

$$E_e \text{ and } b_s \text{ are constant} \quad N_1(x') = 1 - \frac{x'}{h_e} \quad \text{and} \quad N_2(x') = \frac{x'}{h_e}$$

$$\mathbf{K}^e = \frac{E_e}{h_e} \begin{bmatrix} 1 & -1 \\ -1 & 1 \end{bmatrix} \quad \text{and} \quad \mathbf{f}^e = \frac{1}{2} b_s h_e \begin{Bmatrix} 1 \\ 1 \end{Bmatrix}$$

7

### 3.3 Finite element computation based on weak form

- Weak form in global domain

$$G(w, u) = \int_{\Omega} w(x) b_s dx - \int_{\Omega} \frac{\partial w}{\partial x} E \frac{\partial u}{\partial x} dx + w \bar{t}_x \Big|_{\Gamma_r} = 0$$

- Weak form in each element

$$\hat{G}(\hat{w}, \hat{u}) = \hat{G}_e(\hat{w}, \hat{u}) - \hat{G}_f(\hat{w}, \hat{u}) - \hat{w}(x) \bar{t}_x \Big|_{\Gamma_r} = 0$$

$$\hat{G}_e(\hat{w}, \hat{u}) = \sum_{e=1}^M [\hat{w}_1^e \quad \hat{w}_2^e] \int_0^{h_e} \begin{Bmatrix} \frac{dN_1}{dx'} \\ \frac{dN_2}{dx'} \end{Bmatrix} E_s \begin{bmatrix} \frac{dN_1}{dx'} & \frac{dN_2}{dx'} \end{bmatrix} dx' \begin{Bmatrix} \hat{u}_1^e \\ \hat{u}_2^e \end{Bmatrix} = \sum_{e=1}^M [\hat{w}_1^e \quad \hat{w}_2^e] \mathbf{K}^e \begin{Bmatrix} \hat{u}_1^e \\ \hat{u}_2^e \end{Bmatrix}$$

$$\hat{G}_f(\hat{w}, \hat{u}) = \sum_{e=1}^M [\hat{w}_1^e \quad \hat{w}_2^e] \int_0^{h_e} \begin{Bmatrix} N_1 \\ N_2 \end{Bmatrix} b_s dx' = \sum_{e=1}^M [\hat{w}_1^e \quad \hat{w}_2^e] \mathbf{f}^e dx'$$

8

### 3.3 Finite element computation based on weak form

- Weak form in global domain

$$G(w, u) = \int_{\Omega} w(x) b_s dx - \int_{\Omega} \frac{\partial w}{\partial x} E \frac{\partial u}{\partial x} dx + w \bar{t}_x \Big|_{\Gamma_r} = 0$$

- Weak form in each element

$$\hat{G}(\hat{w}, \hat{u}) = \hat{G}_e(\hat{w}, \hat{u}) - \hat{G}_f(\hat{w}, \hat{u}) - \hat{w}(x) \bar{t}_x \Big|_{\Gamma_r} = 0$$

- Since  $\begin{bmatrix} \hat{w}_1^e \\ \hat{w}_2^e \end{bmatrix} = \begin{bmatrix} \hat{w}_1^{e-1} \\ \hat{w}_1^e \end{bmatrix}$

$$\hat{w}_1^1 f_1(\mathbf{K}^e, \mathbf{f}^e, \mathbf{u}^e) + \hat{w}_1^2 f_2(\mathbf{K}^e, \mathbf{f}^e, \mathbf{u}^e) + \dots + \hat{w}_1^{N-1} f_{N-1}(\mathbf{K}^e, \mathbf{f}^e, \mathbf{u}^e) + \hat{w}_2^{N-1} f_N(\mathbf{K}^e, \mathbf{f}^e, \mathbf{u}^e) = 0$$

9

### 3.3 Finite element computation based on weak form

- Weak form in global domain

$$G(w, u) = \int_{\Omega} w(x) b_s dx - \int_{\Omega} \frac{\partial w}{\partial x} E \frac{\partial u}{\partial x} dx + w \bar{t}_x \Big|_{\Gamma_r} = 0$$

- Weak form in each element

$$\hat{G}(\hat{w}, \hat{u}) = \hat{G}_e(\hat{w}, \hat{u}) - \hat{G}_f(\hat{w}, \hat{u}) - \hat{w}(x) \bar{t}_x \Big|_{\Gamma_r} = 0$$

- Since the parameters  $\hat{w}_i^e$  are arbitrary, the expression multiplying each one must be zero. This leads to the set of equations:

$$\Rightarrow \begin{cases} f_1(\mathbf{K}^e, \mathbf{f}^e, \mathbf{u}^e) = 0 \\ f_2(\mathbf{K}^e, \mathbf{f}^e, \mathbf{u}^e) = 0 \\ \vdots \\ f_{N-1}(\mathbf{K}^e, \mathbf{f}^e, \mathbf{u}^e) = 0 \\ f_N(\mathbf{K}^e, \mathbf{f}^e, \mathbf{u}^e) = 0 \end{cases} \Rightarrow \mathbf{K}\mathbf{u} - \mathbf{f} = \mathbf{0} \Rightarrow \mathbf{K}\mathbf{u} = \mathbf{f}$$

10

# The End