

Computational Mechanics

(English Course)

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3 Weak Form of Equivalent Integration



- 3.1 Weak form of equivalent integration for differential equations
- 3.2 Weak form of one-dimensional elasticity problems
- 3.3 Finite element computation based on weak form
- 3.4 Global assembly from one-dimensional elements
- 3.5 Treatments on boundary conditions
- 3.6 Exercises

3.3 Finite element computation based on weak form



Finite element computation

- A more convenient method to construct the approximating functions ϕ_n and ψ_m are obtained by dividing the domain to be analyzed into small regular shaped regions. For example, we can divide the one-dimensional region between a and b into a set of “ M ” finite small segments by defining a set of N points x_i such that

$$x_1 = a, \quad x_i < x_{i+1} \quad \text{and} \quad x_N = b$$

For a one-dimensional problem we can let each increment define a *finite element* domain (or more simply, an *element*) and the set of points define the *nodes* (*finite element mesh* or *mesh*).

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3.3 Finite element computation based on weak form



- A simple set of continuous polynomial approximating functions

$$\phi_i = \begin{cases} 0, & x < x_{i-1} \\ \frac{x - x_{i-1}}{x_i - x_{i-1}}, & x_{i-1} \leq x \leq x_i \\ \frac{x_{i+1} - x}{x_{i+1} - x_i}, & x_i < x \leq x_{i+1} \\ 0, & x > x_{i+1} \end{cases}$$

C_0 function since only the function is continuous in x , whereas the first derivative is only piecewise continuous with the discontinuities located at the nodes.

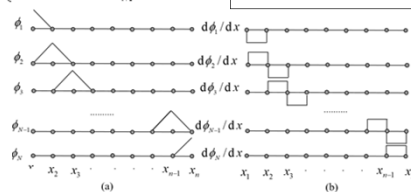


Figure 3.3 One-dimensional finite element approximation for ϕ_i : (a) functions and (b) derivatives.

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3.3 Finite element computation based on weak form



- Integrals over each element in the weak form

$$\int_{\Omega} (\cdot) dx = \sum_{i=1}^M \int_{x_{i-1}}^{x_i} (\cdot) dx \equiv \sum_e \int_{\Omega_e} (\cdot) dx$$

Considering any interval $[x_i, x_{i+1}]$, we note that each interval is defined by the same two local functions N_1 and N_2 . We call these the *shape functions* (形函数) for the element.

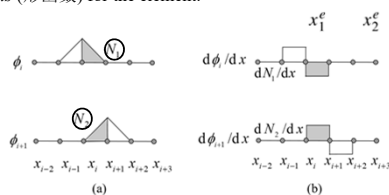


Figure 3.4 One-dimensional finite element shape functions: (a) functions and (b) derivatives.

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3.3 Finite element computation based on weak form



- Define local nodal coordinates on each element as x_1^e and x_2^e .

$$x_1^e \quad \left| \begin{array}{c} e \\ \hline x' \\ \hline \end{array} \right| \quad x_2^e$$

$$x' = x - x_1^e \quad h_e = x_2^e - x_1^e$$

- Local coordinate system

$$N_1(x') = 1 - \frac{x'}{h_e} \quad \text{and} \quad N_2(x') = \frac{x'}{h_e}$$

$$\frac{dN_1}{dx} = \frac{dN_1}{dx'} = -\frac{1}{h_e} \quad \text{and} \quad \frac{dN_2}{dx} = \frac{dN_2}{dx'} = \frac{1}{h_e}$$

- Displacements and arbitrary function on the element

$$\hat{u}^e = N_1(x') \hat{u}_1^e + N_2(x') \hat{u}_2^e$$

$$\hat{w}^e = N_1(x') \hat{w}_1^e + N_2(x') \hat{w}_2^e$$

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3.3 Finite element computation based on weak form

- Weak form in global domain

$$G(w, u) = \int_{\Omega} w(x) b_x dx - \int_{\Omega} \frac{\partial w}{\partial x} E \frac{\partial u}{\partial x} dx + w \bar{t}_x \Big|_{\Gamma_r} = 0$$

- Weak form in each element

$$\hat{G}(\hat{w}, \hat{u}) = \hat{G}_e(\hat{w}, \hat{u}) - \hat{G}_f(\hat{w}, \hat{u}) - \hat{w}(x) \bar{t}_x \Big|_{\Gamma_r} = 0$$

$$\hat{G}_e(\hat{w}, \hat{u}) = \sum_{e=1}^M [\hat{w}_1^e \quad \hat{w}_2^e] \int_0^{h_e} \begin{Bmatrix} \frac{dN_1}{dx'} \\ \frac{dN_2}{dx'} \end{Bmatrix} E \begin{bmatrix} \frac{dN_1}{dx'} & \frac{dN_2}{dx'} \end{bmatrix} dx' \begin{Bmatrix} \hat{u}_1^e \\ \hat{u}_2^e \end{Bmatrix}$$

$$\hat{G}_f(\hat{w}, \hat{u}) = \sum_{e=1}^M [\hat{w}_1^e \quad \hat{w}_2^e] \int_0^{h_e} \begin{Bmatrix} N_1 \\ N_2 \end{Bmatrix} b_x dx'$$

3.3 Finite element computation based on weak form

- Each element can be evaluated as element stiffness matrix and load matrix

$$\mathbf{K}^e = \int_0^{h_e} \begin{Bmatrix} \frac{dN_1}{dx'} \\ \frac{dN_2}{dx'} \end{Bmatrix} E \begin{bmatrix} \frac{dN_1}{dx'} & \frac{dN_2}{dx'} \end{bmatrix} dx' = \begin{bmatrix} K_{11}^e & K_{12}^e \\ K_{21}^e & K_{22}^e \end{bmatrix}$$

$$\mathbf{f}^e = \int_0^{h_e} \begin{Bmatrix} N_1 \\ N_2 \end{Bmatrix} b_x dx' = \begin{Bmatrix} f_1^e \\ f_2^e \end{Bmatrix}$$

- Element stiffness matrix and load matrix

$$E_e \text{ and } b_x \text{ are constant} \quad N_1(x') = 1 - \frac{x'}{h_e} \quad \text{and} \quad N_2(x') = \frac{x'}{h_e}$$

$$\mathbf{K}^e = \frac{E_e}{h_e} \begin{bmatrix} 1 & -1 \\ -1 & 1 \end{bmatrix} \quad \text{and} \quad \mathbf{f}^e = \frac{1}{2} b_x h_e \begin{Bmatrix} 1 \\ 1 \end{Bmatrix}$$

3.3 Finite element computation based on weak form

- Weak form in global domain

$$G(w, u) = \int_{\Omega} w(x) b_x dx - \int_{\Omega} \frac{\partial w}{\partial x} E \frac{\partial u}{\partial x} dx + w \bar{t}_x \Big|_{\Gamma_r} = 0$$

- Weak form in each element

$$\hat{G}(\hat{w}, \hat{u}) = \hat{G}_e(\hat{w}, \hat{u}) - \hat{G}_f(\hat{w}, \hat{u}) - \hat{w}(x) \bar{t}_x \Big|_{\Gamma_r} = 0$$

$$\hat{G}_e(\hat{w}, \hat{u}) = \sum_{e=1}^M [\hat{w}_1^e \quad \hat{w}_2^e] \int_0^{h_e} \begin{Bmatrix} \frac{dN_1}{dx'} \\ \frac{dN_2}{dx'} \end{Bmatrix} E_e \begin{bmatrix} \frac{dN_1}{dx'} & \frac{dN_2}{dx'} \end{bmatrix} dx' \begin{Bmatrix} \hat{u}_1^e \\ \hat{u}_2^e \end{Bmatrix} = \sum_{e=1}^M [\hat{w}_1^e \quad \hat{w}_2^e] \mathbf{K}^e \begin{Bmatrix} \hat{u}_1^e \\ \hat{u}_2^e \end{Bmatrix}$$

$$\hat{G}_f(\hat{w}, \hat{u}) = \sum_{e=1}^M [\hat{w}_1^e \quad \hat{w}_2^e] \int_0^{h_e} \begin{Bmatrix} N_1 \\ N_2 \end{Bmatrix} b_x dx' = \sum_{e=1}^M [\hat{w}_1^e \quad \hat{w}_2^e] \mathbf{f}^e dx'$$

3.3 Finite element computation based on weak form

- Weak form in global domain

$$G(w, u) = \int_{\Omega} w(x) b_x dx - \int_{\Omega} \frac{\partial w}{\partial x} E \frac{\partial u}{\partial x} dx + w \bar{t}_x \Big|_{\Gamma_r} = 0$$

- Weak form in each element

$$\hat{G}(\hat{w}, \hat{u}) = \hat{G}_e(\hat{w}, \hat{u}) - \hat{G}_f(\hat{w}, \hat{u}) - \hat{w}(x) \bar{t}_x \Big|_{\Gamma_r} = 0$$

- Since $[\hat{w}_1^e \quad \hat{w}_2^e] = \underline{\hat{w}}_e^T$

$$\hat{w}_1^e f_1^e (\mathbf{K}^e, \mathbf{f}^e, \mathbf{u}^e) + \hat{w}_2^e f_2^e (\mathbf{K}^e, \mathbf{f}^e, \mathbf{u}^e) + \dots + \hat{w}_1^{e-1} f_{N-1}^e (\mathbf{K}^e, \mathbf{f}^e, \mathbf{u}^e) + \hat{w}_2^{e-1} f_N^e (\mathbf{K}^e, \mathbf{f}^e, \mathbf{u}^e) = 0$$

3.3 Finite element computation based on weak form

- Weak form in global domain

$$G(w, u) = \int_{\Omega} w(x) b_x dx - \int_{\Omega} \frac{\partial w}{\partial x} E \frac{\partial u}{\partial x} dx + w \bar{t}_x \Big|_{\Gamma_r} = 0$$

- Weak form in each element

$$\hat{G}(\hat{w}, \hat{u}) = \hat{G}_e(\hat{w}, \hat{u}) - \hat{G}_f(\hat{w}, \hat{u}) - \hat{w}(x) \bar{t}_x \Big|_{\Gamma_r} = 0$$

- Since the parameters \hat{w}_i^e are arbitrary, the expression multiplying each one must be zero. This leads to the set of equations:

$$\Rightarrow \begin{cases} f_1(\mathbf{K}^e, \mathbf{f}^e, \mathbf{u}^e) = 0 \\ f_2(\mathbf{K}^e, \mathbf{f}^e, \mathbf{u}^e) = 0 \\ \vdots \\ f_{N-1}(\mathbf{K}^e, \mathbf{f}^e, \mathbf{u}^e) = 0 \\ f_N(\mathbf{K}^e, \mathbf{f}^e, \mathbf{u}^e) = 0 \end{cases} \Rightarrow \mathbf{Ku} - \mathbf{f} = \mathbf{0} \Rightarrow \mathbf{Ku} = \mathbf{f}$$

3 Weak Form of Equivalent Integration

3.1 Weak form of equivalent integration for differential equations

3.2 Weak form of one-dimensional elasticity problems

3.3 Finite element computation based on weak form

3.4 Global assembly from one-dimensional elements

3.5 Treatments on boundary conditions

3.6 Exercises

3.4 Global assembly from one-dimensional elements

- In the one-dimensional beam model, $N-1$ elements ("M" small finite segments) and N nodes are used



Figure 3.5 Element and node numbers of one-dimensional problem.

3.4 Global assembly from one-dimensional elements

- Then for each element we define the relationship of the local nodes to the global node number

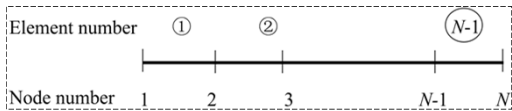
Table 3.2 Local to global node numbering for two-end element for one-dimensional elastic beam.

Local node number	Element number					
	1	2	3	...	N-1	
1	1	2	3	...	N-1	
2	2	3	4	...	N	

3.4 Global assembly from one-dimensional elements

- According to the entire node number of each element, the corresponding element location vector can be provided for determining the relationships between local and global node locations as follows

$$\lambda^1 = \{1 \ 2\}^T \quad \lambda^2 = \{2 \ 3\}^T \quad \dots \quad \lambda^{N-1} = \{N-1 \ N\}^T$$



3.4 Global assembly from one-dimensional elements

- In order to implement global assembly from one-dimensional elements, the node numbers are marked on the left and upper sides of the matrix, and the expanded form of the stiffness matrix is given via element location vector

$$\begin{matrix} \text{Node} \\ \text{number} \end{matrix} \begin{matrix} \leftarrow & 1 & 2 & 3 & \dots & N-1 & N \\ \begin{matrix} 1 \\ 2 \\ 3 \\ \vdots \\ N-1 \\ N \end{matrix} & \begin{bmatrix} K_{11}^1 & K_{12}^1 & 0 & & & \\ K_{21}^1 & K_{22}^1 + K_{11}^2 & K_{12}^2 & & & \\ 0 & K_{21}^2 & (K_{22}^2 + K_{11}^3) & & & \\ \vdots & & & \ddots & & \\ K_{N-1,1}^{N-1} & & & & (K_{22}^{N-1} + K_{11}^N) & K_{12}^N \\ & & & & K_{21}^N & K_{22}^N \end{bmatrix} & \begin{bmatrix} \hat{u}_1 \\ \hat{u}_2 \\ \hat{u}_3 \\ \vdots \\ \hat{u}_{N-1} \\ \hat{u}_N \end{bmatrix} \end{matrix} = \begin{bmatrix} f_1 \\ f_2 \\ f_3 \\ \vdots \\ f_{N-1} \\ f_N \end{bmatrix}$$

3.4 Global assembly from one-dimensional elements

- Similarly, the of node numbers is marked on the left of the load vector, and the expanded load is given by

$$\begin{matrix} \text{Node} \\ \text{number} \end{matrix} \begin{matrix} \downarrow \\ 1 \\ 2 \\ 3 \\ \vdots \\ N-1 \\ N \end{matrix} \begin{bmatrix} f_1 \\ f_2 \\ f_3 \\ \vdots \\ f_{N-1} \\ f_N \end{bmatrix} = \begin{bmatrix} f_1^1 \\ f_2^1 + f_1^2 \\ f_2^2 + f_1^3 \\ \vdots \\ f_2^{N-1} + f_1^N \\ f_2^N \end{bmatrix}$$

3.4 Global assembly from one-dimensional elements

- A standard linear problem with the final stiffness equations is as follows

$$Ku = f$$

$$\begin{matrix} \text{Node} \\ \text{number} \end{matrix} \begin{matrix} \leftarrow & 1 & 2 & 3 & \dots & N-1 & N \\ \begin{matrix} 1 \\ 2 \\ 3 \\ \vdots \\ N-1 \\ N \end{matrix} & \begin{bmatrix} K_{11}^1 & K_{12}^1 & 0 & & & \\ K_{21}^1 & K_{22}^1 + K_{11}^2 & K_{12}^2 & & & \\ 0 & K_{21}^2 & (K_{22}^2 + K_{11}^3) & & & \\ \vdots & & & \ddots & & \\ K_{N-1,1}^{N-1} & & & & (K_{22}^{N-1} + K_{11}^N) & K_{12}^N \\ & & & & K_{21}^N & K_{22}^N \end{bmatrix} & \begin{bmatrix} \hat{u}_1 \\ \hat{u}_2 \\ \hat{u}_3 \\ \vdots \\ \hat{u}_{N-1} \\ \hat{u}_N \end{bmatrix} \end{matrix} = \begin{bmatrix} f_1 \\ f_2 \\ f_3 \\ \vdots \\ f_{N-1} \\ f_N \end{bmatrix}$$

3.4 Global assembly from one-dimensional elements

- Matrix form of stiffness equation

$$\mathbf{Ku} = \mathbf{f}$$

$$\mathbf{K} = \begin{bmatrix} K_{11} & K_{12} & \dots & K_{1N} \\ K_{21} & K_{22} & & K_{2N} \\ \vdots & & & \vdots \\ K_{N1} & K_{N2} & \dots & K_{NN} \end{bmatrix} \text{ and } \mathbf{f} = \begin{bmatrix} f_1 \\ f_2 \\ \vdots \\ f_N \end{bmatrix}$$

- Formal solution

$$\mathbf{u} = \mathbf{K}^{-1}\mathbf{f}$$

- Displacement and stress solution

$$\hat{u}'(x) = N_1(x)\hat{u}'_1 + N_2(x)\hat{u}'_2 \quad \hat{\sigma}_x(x) = E \frac{\partial \hat{u}(x)}{\partial x} = E \left(\frac{dN_1(x)}{dx} \hat{u}'_1 + \frac{dN_2(x)}{dx} \hat{u}'_2 \right)$$



3 Weak Form of Equivalent Integration



3.1 Weak form of equivalent integration for differential equations

3.2 Weak form of one-dimensional elasticity problems

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3.6 Exercises

3.5 Treatments on boundary conditions



- (1) Displacement boundary conditions $\hat{u}(x_a) \equiv \bar{u}_a \quad \hat{u}_1 \equiv \bar{u}_1$

$$\begin{bmatrix} 1 & 0 & 0 \\ 0 & K_{22}^1 + K_{11}^2 & K_{12}^2 & 0 \\ 0 & K_{21}^2 & (K_{22}^2 + K_{11}^3) & K_{12}^3 \\ & & & \ddots \\ & & K_{21}^{N-1} & (K_{22}^{N-1} + K_{11}^N) & K_{12}^N \\ & & & K_{21}^N & K_{22}^N \end{bmatrix} \begin{bmatrix} \hat{u}_1 \\ \hat{u}_2 \\ \hat{u}_3 \\ \vdots \\ \hat{u}_{N-1} \\ \hat{u}_N \end{bmatrix} = \begin{bmatrix} \bar{u}_1 \\ f_2 - K_{21}^1 \bar{u}_1 \\ f_3 \\ \vdots \\ f_{N-1} \\ f_N \end{bmatrix}$$

- (2) Traction boundary conditions $x = L$

$$f_N \rightarrow f_N + \bar{t}_x(L)$$

3.3 Finite element computation based on weak form



- Weak form in global domain

$$G(w, u) = \int_a^b w(x) b_x dx - \int_a^b \frac{\partial w}{\partial x} E \frac{\partial u}{\partial x} dx + w \bar{t}_x \Big|_r = 0$$

- Weak form in each element

$$\hat{G}(\hat{w}, \hat{u}) = \hat{G}_e(\hat{w}, \hat{u}) - \hat{G}_f(\hat{w}, \hat{u}) - \hat{w}(x) \bar{t}_x \Big|_r$$

$$\hat{G}_e(\hat{w}, \hat{u}) = \sum_{e=1}^M \int_0^{h_e} \left[\hat{w}_1^e \hat{w}_2^e \right]_0^{h_e} E \left[\frac{dN_1}{dx'} \frac{dN_2}{dx'} \right] dx' \begin{bmatrix} \hat{u}_1^e \\ \hat{u}_2^e \end{bmatrix}$$

$$\hat{G}_f(\hat{w}, \hat{u}) = \sum_{e=1}^M \int_0^{h_e} \left[\hat{w}_1^e \quad \hat{w}_2^e \right]_0^{h_e} \begin{bmatrix} N_1 \\ N_2 \end{bmatrix} b_x dx'$$

Review previous content

3.5 Treatments on boundary conditions



Example 3.2. Solutions of finite element method for one-dimensional elastic cantilever beam

- As an example we consider a static problem with length 10 units and $E=1000$. Figure 3.1 shows the problem to solve.

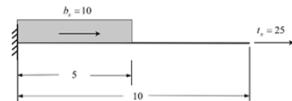


Figure 3.1 One-dimensional elastic cantilever beam in Example 3.1.

- Loading $b_x = \begin{cases} 10 & \text{for } 0 < x < 5 \\ 0 & \text{for } 5 < x < 10 \end{cases}$

- Boundary conditions $u(0) = 0$ and $\bar{t}_x(10) = 25$

- We divide the domain into four equal elements.

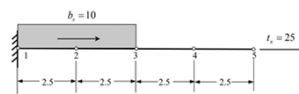


Figure 3.6 Four-element mesh for one-dimensional elastic cantilever beam in Example 3.2.

3.3 Finite element computation based on weak form



- Each element can be evaluated as element stiffness matrix and load matrix

$$\mathbf{K}^e = \int_0^{h_e} \begin{bmatrix} \frac{dN_1}{dx'} \\ \frac{dN_2}{dx'} \end{bmatrix} E \begin{bmatrix} \frac{dN_1}{dx'} & \frac{dN_2}{dx'} \end{bmatrix} dx' = \begin{bmatrix} K_{11}^e & K_{12}^e \\ K_{21}^e & K_{22}^e \end{bmatrix}$$

$$\mathbf{f}^e = \int_0^{h_e} \begin{bmatrix} N_1 \\ N_2 \end{bmatrix} b_x dx' = \begin{bmatrix} f_1^e \\ f_2^e \end{bmatrix}$$

Review previous content

- Element stiffness matrix and load matrix

$$E_e \text{ and } b_x \text{ are constant} \quad N_1(x') = 1 - \frac{x'}{h_e} \quad \text{and} \quad N_2(x') = \frac{x'}{h_e}$$

$$\mathbf{K}^e = \frac{E_e}{h_e} \begin{bmatrix} 1 & -1 \\ -1 & 1 \end{bmatrix} \quad \text{and} \quad \mathbf{f}^e = \frac{1}{2} b_x h_e \begin{bmatrix} 1 \\ 1 \end{bmatrix}$$

3.5 Treatments on boundary conditions

- Element stiffness matrix and load matrix

$$\mathbf{K}^1 = \mathbf{K}^2 = \mathbf{K}^3 = \mathbf{K}^4 = \frac{E_e}{h_e} \begin{bmatrix} 1 & -1 \\ -1 & 1 \end{bmatrix} = \frac{1000}{2.5} \begin{bmatrix} 1 & -1 \\ -1 & 1 \end{bmatrix} = \begin{bmatrix} 400 & -400 \\ -400 & 400 \end{bmatrix}$$

$$\mathbf{f}^1 = \mathbf{f}^2 = \frac{1}{2} b_1 h_e \begin{Bmatrix} 1 \\ 1 \end{Bmatrix} = \frac{1}{2} \times 10 \times 2.5 \begin{Bmatrix} 1 \\ 1 \end{Bmatrix} = \begin{Bmatrix} 12.5 \\ 12.5 \end{Bmatrix}$$

$$\mathbf{f}^3 = \mathbf{f}^4 = \frac{1}{2} b_1 h_e \begin{Bmatrix} 1 \\ 1 \end{Bmatrix} = \frac{1}{2} \times 0 \times 2.5 \begin{Bmatrix} 1 \\ 1 \end{Bmatrix} = \begin{Bmatrix} 0 \\ 0 \end{Bmatrix}$$

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3.5 Treatments on boundary conditions

- Element stiffness matrix

Node number	1	2	3	4	5
	K_{11}^1	K_{12}^1	K_{13}^2	K_{14}^3	K_{15}^4
	K_{21}^1	$(K_{22}^1 + K_{22}^2)$	K_{23}^2	K_{24}^3	K_{25}^4
		K_{31}^2	$(K_{32}^2 + K_{32}^3)$	K_{34}^3	K_{35}^4
			K_{41}^3	$(K_{42}^3 + K_{42}^4)$	K_{45}^4
				K_{51}^4	K_{52}^4

$$= \begin{bmatrix} 400 & -400 & & & \\ -400 & 400 + 400 & -400 & & \\ & -400 & 400 + 400 & -400 & \\ & & -400 & 400 + 400 & -400 \\ & & & -400 & 400 \end{bmatrix}$$

$$= \begin{bmatrix} 400 & -400 & & & \\ -400 & 800 & -400 & & \\ & -400 & 800 & -400 & \\ & & -400 & 800 & -400 \\ & & & -400 & 400 \end{bmatrix}$$

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3.5 Treatments on boundary conditions

- Load matrix

Node number	1	2	3	4	5
	f_1	f_2	f_3	f_4	f_5
	f_1^1	$f_2^1 + f_2^2$	$f_3^2 + f_3^3$	$f_4^3 + f_4^4$	f_5^4
	12.5	12.5 + 12.5	12.5 + 0	0 + 0	0
	12.5	25	12.5	0	0

- Matrix form of stiffness equation $\mathbf{Ku} = \mathbf{f}$

$$\begin{bmatrix} 400 & -400 & & & \\ -400 & 800 & -400 & & \\ & -400 & 800 & -400 & \\ & & -400 & 800 & -400 \\ & & & -400 & 400 \end{bmatrix} \begin{Bmatrix} \hat{u}_1 \\ \hat{u}_2 \\ \hat{u}_3 \\ \hat{u}_4 \\ \hat{u}_5 \end{Bmatrix} = \begin{Bmatrix} 12.5 \\ 25 \\ 12.5 \\ 0 \\ 0 \end{Bmatrix}$$

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3.5 Treatments on boundary conditions

- Displacement boundary conditions $\hat{u}_1 \equiv \bar{u}_1 = 0$

$$\begin{bmatrix} 400 & -400 & & & \\ -400 & 800 & -400 & & \\ & -400 & 800 & -400 & \\ & & -400 & 800 & -400 \\ & & & -400 & 400 \end{bmatrix} \begin{Bmatrix} \hat{u}_1 \\ \hat{u}_2 \\ \hat{u}_3 \\ \hat{u}_4 \\ \hat{u}_5 \end{Bmatrix} = \begin{Bmatrix} 0 \\ 25 \\ 12.5 \\ 0 \\ 0 \end{Bmatrix}$$

$$\Rightarrow \begin{bmatrix} 0 & 800 & -400 & & \\ -400 & 800 & -400 & & \\ & -400 & 800 & -400 & \\ & & -400 & 800 & -400 \\ & & & -400 & 400 \end{bmatrix} \begin{Bmatrix} \hat{u}_1 \\ \hat{u}_2 \\ \hat{u}_3 \\ \hat{u}_4 \\ \hat{u}_5 \end{Bmatrix} = \begin{Bmatrix} 0 \\ 25 \\ 12.5 \\ 0 \\ 0 \end{Bmatrix} \quad f_2 - K_{21}^1 \bar{u}_1$$

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3.5 Treatments on boundary conditions

- Traction boundary conditions $f_5 \rightarrow f_5 + \bar{f}_5(10) = 25$

$$\begin{bmatrix} 800 & -400 & & & \\ -400 & 800 & -400 & & \\ & -400 & 800 & -400 & \\ & & -400 & 800 & -400 \\ & & & -400 & 400 \end{bmatrix} \begin{Bmatrix} \hat{u}_1 \\ \hat{u}_2 \\ \hat{u}_3 \\ \hat{u}_4 \\ \hat{u}_5 \end{Bmatrix} = \begin{Bmatrix} 0 \\ 25 \\ 12.5 \\ 0 \\ 0 \end{Bmatrix}$$

$$\Rightarrow \begin{bmatrix} 800 & -400 & & & \\ -400 & 800 & -400 & & \\ & -400 & 800 & -400 & \\ & & -400 & 800 & -400 \\ & & & -400 & 400 \end{bmatrix} \begin{Bmatrix} \hat{u}_1 \\ \hat{u}_2 \\ \hat{u}_3 \\ \hat{u}_4 \\ \hat{u}_5 \end{Bmatrix} = \begin{Bmatrix} 0 \\ 25 \\ 12.5 \\ 0 \\ 25 \end{Bmatrix}$$

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3.5 Treatments on boundary conditions

- Formal solution $\mathbf{u} = \mathbf{K}^{-1} \mathbf{f}$

$$\begin{bmatrix} 800 & -400 & & & \\ -400 & 800 & -400 & & \\ & -400 & 800 & -400 & \\ & & -400 & 800 & -400 \\ & & & -400 & 400 \end{bmatrix} \begin{Bmatrix} \hat{u}_1 \\ \hat{u}_2 \\ \hat{u}_3 \\ \hat{u}_4 \\ \hat{u}_5 \end{Bmatrix} = \begin{Bmatrix} 0 \\ 25 \\ 12.5 \\ 0 \\ 25 \end{Bmatrix} \Rightarrow \begin{Bmatrix} \hat{u}_1 \\ \hat{u}_2 \\ \hat{u}_3 \\ \hat{u}_4 \\ \hat{u}_5 \end{Bmatrix} = \begin{Bmatrix} 0 \\ 0.15625 \\ 0.25 \\ 0.3125 \\ 0.375 \end{Bmatrix}$$

- Displacement and stress solution

$$\hat{u}^i(x') = N_1(x') \hat{u}_1^e + N_2(x') \hat{u}_2^e \quad \hat{\sigma}_x^i(x') = E \frac{\partial \hat{u}^i(x')}{\partial x'} = E \left(\frac{dN_1(x')}{dx'} \hat{u}_1^e + \frac{dN_2(x')}{dx'} \hat{u}_2^e \right)$$

$$\hat{u}^i(x') = \left(1 - \frac{x'}{h_e} \right) \times 0 + \frac{x'}{h_e} \times 0.15625 = 0.0625x'$$

$$\hat{\sigma}_x^i(x') = E \frac{\partial \hat{u}^i(x')}{\partial x'} = E \left(-\frac{1}{h_e} \times 0 + \frac{1}{h_e} \times 0.15625 \right) = 62.5$$

62

3.5 Treatments on boundary conditions



- Displacement and stress solution

$$\hat{u}^T(x') = N_1(x')\hat{u}_1^T + N_2(x')\hat{u}_2^T \quad \hat{\sigma}_1^T(x') = E \frac{\partial \hat{u}^T(x')}{\partial x'} = E \left(\frac{dN_1(x')}{dx'}\hat{u}_1^T + \frac{dN_2(x')}{dx'}\hat{u}_2^T \right)$$

$$\hat{u}^1(x') = \left(1 - \frac{x'}{h_e}\right) \times 0.15625 + \frac{x'}{h_e} \times 0.25 = 0.0375x' + 0.15625$$

$$\hat{\sigma}_1^1(x') = E \frac{\partial \hat{u}^1(x')}{\partial x'} = E \left(-\frac{1}{h_e} \times 0.15625 + \frac{1}{h_e} \times 0.25 \right) = 37.5$$

$$\hat{u}^2(x') = \left(1 - \frac{x'}{h_e}\right) \times 0.25 + \frac{x'}{h_e} \times 0.3125 = 0.025x' + 0.25$$

$$\hat{\sigma}_1^2(x') = E \frac{\partial \hat{u}^2(x')}{\partial x'} = E \left(-\frac{1}{h_e} \times 0.25 + \frac{1}{h_e} \times 0.3125 \right) = 25$$

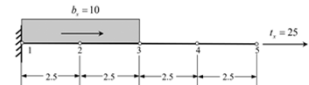
$$\hat{u}^3(x') = \left(1 - \frac{x'}{h_e}\right) \times 0.3125 + \frac{x'}{h_e} \times 0.375 = 0.025x' + 0.3125$$

$$\hat{\sigma}_1^3(x') = E \frac{\partial \hat{u}^3(x')}{\partial x'} = E \left(-\frac{1}{h_e} \times 0.3125 + \frac{1}{h_e} \times 0.375 \right) = 25$$

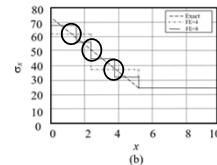
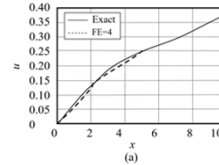
3.5 Treatments on boundary conditions



- We divide the domain into four equal elements.
- Stresses are converging more slowly than the displacements (this was true for the first example also) but again have *super-convergent* points.



• Figure 3.6 Four-element mesh for one-dimensional elastic cantilever beam in Example 3.2.



• Figure 3.7 Displacement and stress solutions in Example 3.2 based on finite element method using N -element solutions: (a) displacement and (b) stress.

Galerkin method VS Finite element method



- Similarities

- 1 Weak form.
- 2 Solution function.
- 3 Algebraic equations.

- Differences

- 1 Domain of solution function: Element, shape function.
- 2 Treatments on boundary conditions: $\mathbf{Ku}=\mathbf{f}$.

The End